

A first order system as a singular limit

of a second order system with low

inertia

In the first few lectures,

we considered the second order equation

$$m\ddot{x} + c\dot{x} + f(x) = 0 \quad - (1)$$

in the limit of $m \rightarrow 0$ (or at least very small m).

We ignored the $m\ddot{x}$ term to obtain

$$c\dot{x} + f(x) = 0$$

$$\dot{x} = \frac{-f(x)}{c} \quad - (2)$$

We noted that (1) required 2 ^{scalar} initial conditions $x(0), \dot{x}(0)$

but (2) requires only $x(0)$.

What is the nature of the relationship between solutions to (1) and (2)?

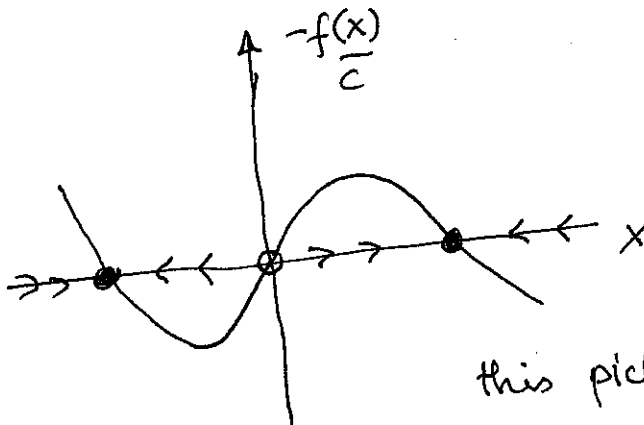
Solutions to (2)

Say

we have

$$\text{this } \frac{-f(x)}{c}$$

with 3 fixed points.



this picture describes what happens for any initial $x(0)$.

Solutions to ①

Equation ① requires 2 scalar initial conditions $x(0)$ and $\dot{x}(0)$. Let us rewrite the equation ① as 2 first order equations:

$$\begin{aligned}\dot{x} &= v && \swarrow \text{velocity.} \\ \dot{v} &= \frac{-(cv + f(x))}{m} && \swarrow \text{acceleration.}\end{aligned}$$

Consider some initial condition $x(0), v(0)$.

$$\dot{v}(0) = \frac{-[cv(0) + f(x(0))]}{m}$$

is HUGE if m is very small

as long as ~~$cv(0) + f(x(0)) \neq 0$~~ $cv(0) + f(x(0)) \neq 0$

$$\text{if } \underline{cv + f(x)} > 0,$$

$$\dot{v} = \text{HUGE and } < 0$$

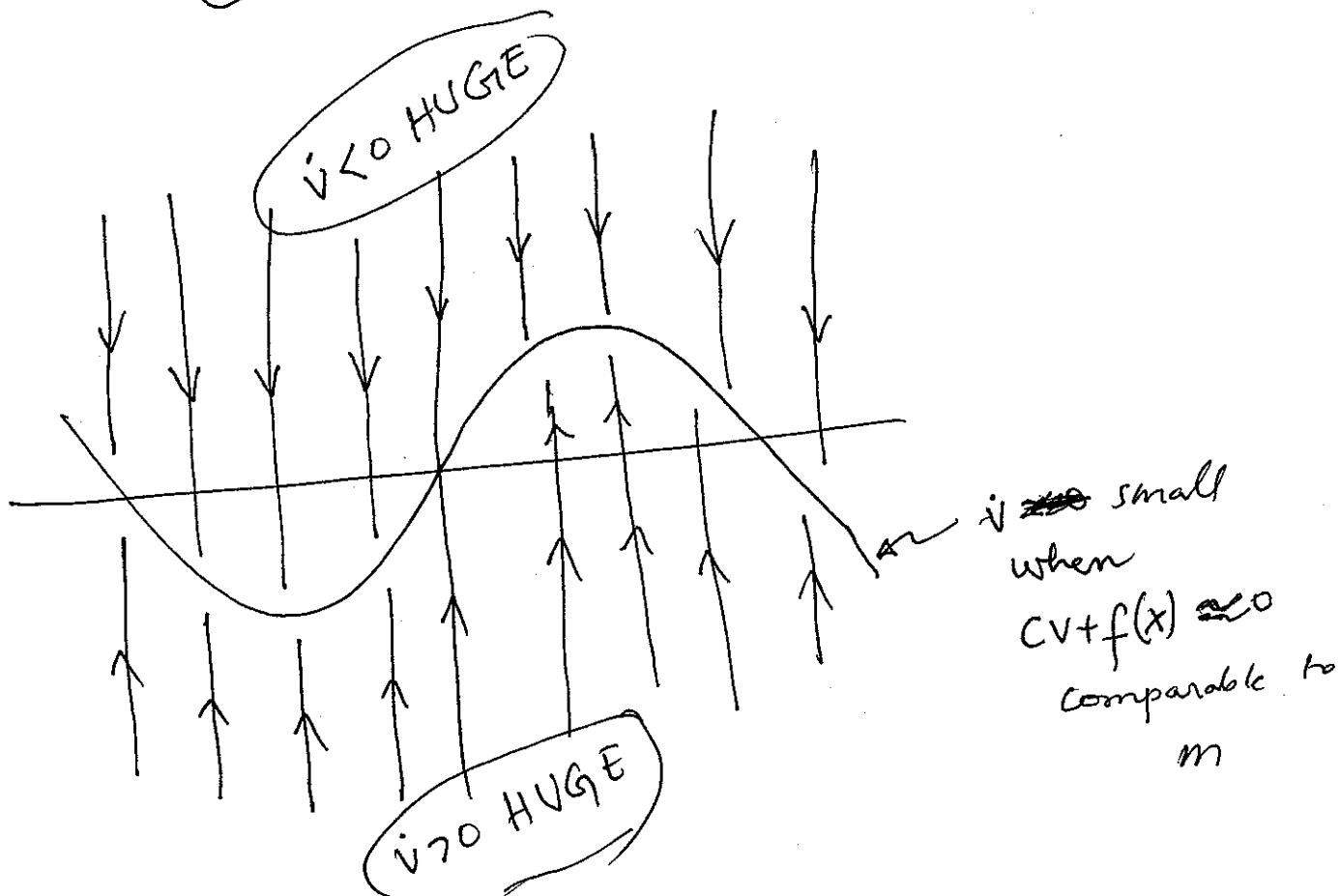
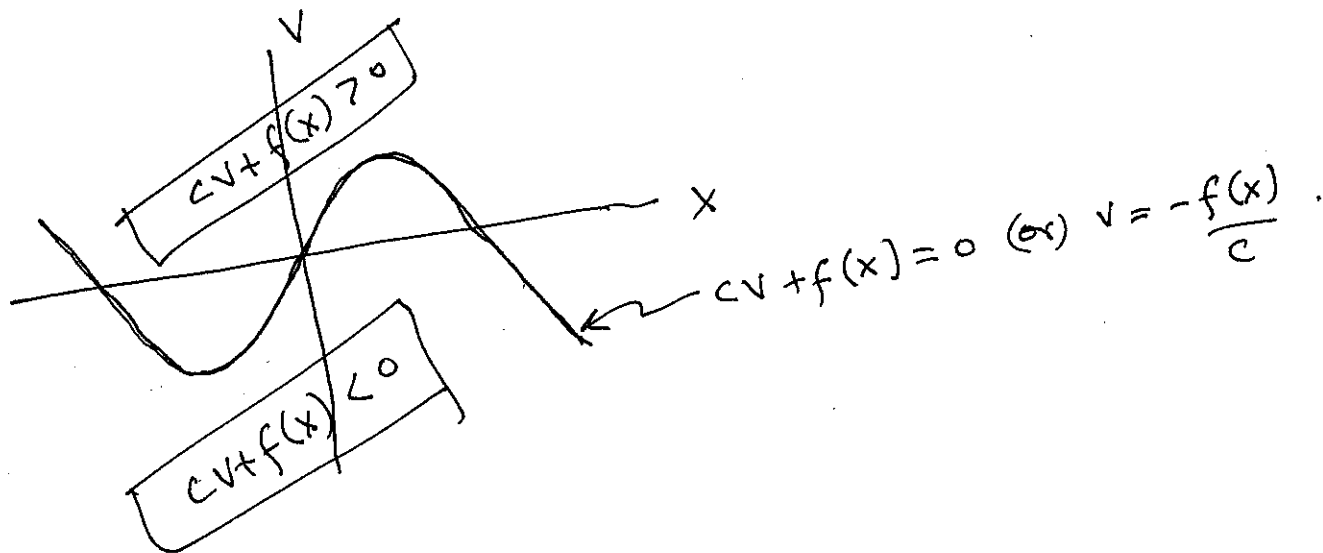
$$\text{if } cv + f(x) < 0$$

$$\dot{v} = \text{HUGE and } > 0.$$

Let's see what this implies in pictures.

first, we divide the phase space into 3 regions.

$$\begin{aligned}cV + f(x) &> 0 \\cV + f(x) &= 0 \\cV + f(x) &< 0.\end{aligned}$$



Thus, when m is VERY small we get HUGE accelerations \ddot{v} that drives the state toward the $cv + f(x)$ curve. VERY FAST.

Once we are ^{very} close to the $cv + f(x)$ curve

\ddot{v} becomes non-HUGE.

and the dynamics are dominated by

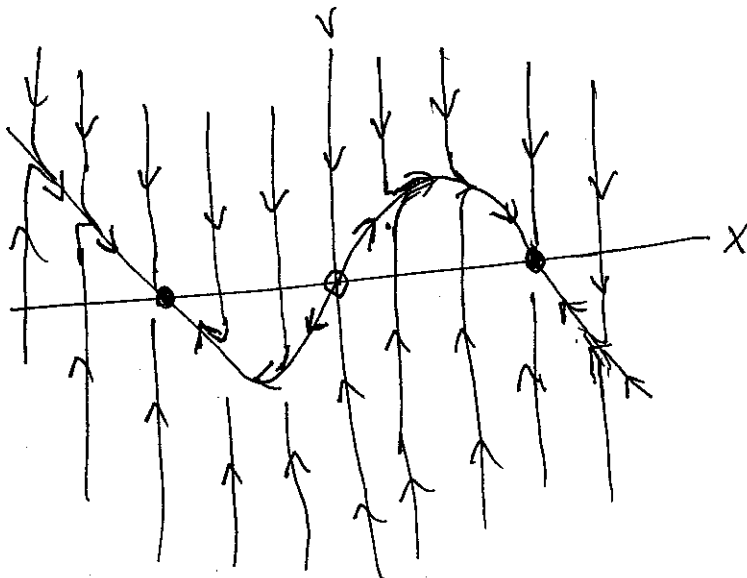
the equation

$$cv + f(x) = 0$$

$$v = -\frac{f(x)}{c} \quad (\text{or})$$

$$\dot{x} = -\frac{f(x)}{c}$$

ie, by the first order ODE. The system crawls along the $cv + f(x) = 0$ toward the appropriate fixed point, shown below.



You can check this phase portrait numerically with $f(x) = x - kx^3$

So this is how the 1st order equation
& the 2nd order equation are related when $m \approx 0$.
When $m \approx 0$, the 2nd order equation has
a solution with 2 pieces

(1) a very fast ^{initial} transient with fast accelerations
that brings the system quickly to

$$cv + f(x) = 0$$

$$(or) v = -f(x)/c.$$

(2) a slow transient that, to a very
good approximation, obeys the
first order ODE

$$\dot{x} = -f(x)/c.$$

The 1st order ~~ODE~~ ODE and the second order ODE
agree except during the brief initial transient.

~ initial brief fast transient is called (sometimes)
a boundary layer (in time)
which is analogous to the spatial boundary
layer in fluid mechanics.

~ this system, having ² time-scales that are
vastly different, [^] are sometimes called

slow-fast systems.

(another example is an earthquake: slow creep over
many years + fast shaking over a few seconds)
we'll see more examples later on.)